

Sivers function and its parameterizations

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Outline

- Parton distributions
- k_\perp dependent parton distributions
- Sivers function
- Interpretation of the Sivers function
- Semi-inclusive deep-inelastic scattering (SIDIS)
- Extraction from HERMES data
- Gauge-Link
- Sivers asymmetry in Drell-Yan (DY)
- Predictions for various experiments
- Summary and Outlook



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Parton distributions

Parton model

Model to describe the distribution of the nucleon's constituents (partons). Mathematical description via quark-quark-correlator

$$\Phi_{ij}(x) = \int \frac{d\xi^-}{4\pi} e^{ixP^+\xi^-} \langle P | \bar{\Psi}_j(0) \mathcal{W}(0, \xi) \Psi_i(\xi^-) | P \rangle$$

collinear parton model

All partons move in the same direction as the nucleon
⇒ no transverse parton momentum

parton distributions

- 1 $f_1^q(x)$ or $q(x)$, e.g. given by $\text{Tr}(\Phi \gamma^+)$
- 2 $g_1^q(x)$ helicity distribution
- 3 $h_1^q(x)$ transversity



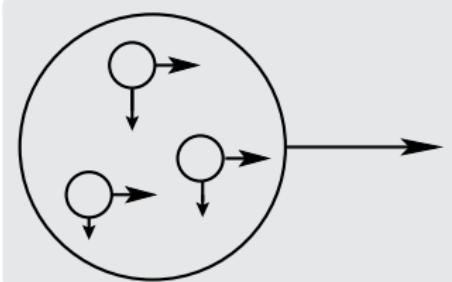
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k_\perp - dependent parton distributions

Parton model with k_\perp

- Partons can also move into transverse direction (compared to the target motion)
- ⇒ more parton distributions, depending on x and also on k_\perp



- Nucleon momentum:
 $P^\mu = (P^+, P^-, \vec{0})$
- Parton momentum:
 $k^\mu = (k^+, k^-, \vec{k}_\perp)$
- By integration over the transverse momentum one gets back to the three forward parton distributions.



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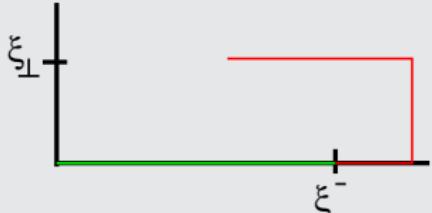
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Quark-Quark correlator

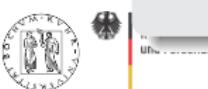
Correlator

$$\Phi_{ij}(x, \vec{k}_\perp) = \int \frac{d\xi^- d^2 \vec{\xi}_\perp}{(2\pi)^3} e^{i(xP^+ \xi^- - \vec{k}_\perp \cdot \vec{\xi}_\perp)} \langle P | \bar{\Psi}_j(0) \mathcal{W}(0, \vec{\xi}) \Psi_i(\vec{\xi}) | P \rangle$$

- $\tilde{\xi} = (0, \xi^-, \vec{\xi}_\perp)$
- $\mathcal{W} = \mathcal{P} \exp(i g \int d\alpha^\mu A_\mu(\alpha))$



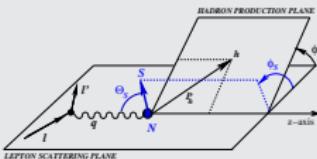
- A_μ : Gluon field
- Light-cone gauge
 $A^+ = 0$
- $\mathcal{W} = 1$ in light-cone gauge, but only if there is no transverse momentum



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Sivers function

- introduced to explain the large Single-Spin asymmetries in e.g. $p^\uparrow p \rightarrow \pi X$ @ FNAL (1990's)
- Measurement of non-zero Sivers Asymmetry also in SIDIS @ HERMES



- $\propto (S_T \times p_T) P_N$
- the Gauge-Link is necessary for the existence of the Sivers effect (e.g. calculation of the Sivers effect in a diquark model)
- the Sivers effect is quantified in terms of the Sivers function f_{1T}^\perp
- Probability to find an unpolarized quark in a transversely polarized target



Sum Rules for Sivers function

Burkardt Sum rule

[M.Burkardt, Phys.Rev. **D69** (2004) 091501]

$$\sum_{a=q,\bar{q},g} \int_0^1 dx \int d^2\vec{k}_\perp \vec{k}_\perp^2 f_{1T}^{\perp a}(x, \vec{k}_\perp^2) = 0$$

$\langle k_\perp \rangle$ vanishes when summed over all partons.

Has been checked in model calculations.

[K. Goeke, S. Meissner, A. Metz, M. Schlegel, Phys.Lett. B637 (2006) 241]



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Connection to GPD's

Connection to $E_q(x, \xi, \Delta)$

[M.Burkardt, Phys.Lett. **B** 595 (2004) 245-249]

- Transform $E_q(x, \xi, \Delta)$ to impact parameter space
 $\mathcal{E}_q(x, b_T)$
- rewrite Sivers function (from scalar diquark spectator model)

- From model calculations

$$E_q \propto (1-x)^5 \quad \Rightarrow \quad f_{1T}^\perp \propto (1-x)^5$$

- Newer calculations predict $f_{1T}^\perp \propto (1-x)^4$
[S.J. Brodsky, F.Yuan, hep-ph/0610236]



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Connection to quark angular momentum

In the calculation of the Sivers function the same light-cone wave functions enter as in the calculation of the anomalous magnetic moment

[S.J. Brodsky, S. Gardner, hep-ph/0608219]

But the SSA matrix elements can only be related to those contributing to κ and are not identical

$$\kappa_p = 1.79$$

$$\kappa_n = -1.91$$

$$\kappa^{u/p} > 0$$

$$\kappa^{d/p} < 0$$

u- and d-quark Sivers
function should have
opposite sign!



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Measuring the Sivers function

- The Sivers function can be measured via a Single-Spin asymmetry (SSA).
- Assume a model for the Sivers function and fit parameters
- Semi-inclusive deep-inelastic scattering off a lepton on a transversely polarized target
- Drell-Yan with one transversely polarized nucleon
- Can also be measured with hadronic final state
 $p^\uparrow p \longrightarrow jet_1 jet_2 X, jet H X$



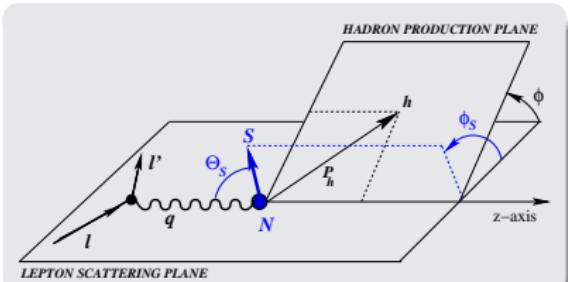
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Semi-inclusive deep-inelastic scattering (SIDIS)

$$A_{UT} \propto \underbrace{\sin(\phi_H + \phi_S) h_1 H_1^\perp}_{\text{Collins Effect}} + \underbrace{\sin(\phi_H - \phi_S) f_{1T}^\perp D_1}_{\text{Sivers Effect}} + \dots$$

$$\sigma = \sigma_{\text{partonic}} \times \text{pdf} \times \text{FF} \times \text{Soft factor}$$



$$q = l - l' \quad Q^2 = -q^2$$

$$x = \frac{Q^2}{2P \cdot q} \quad z = \frac{P \cdot P_H}{P \cdot q}$$

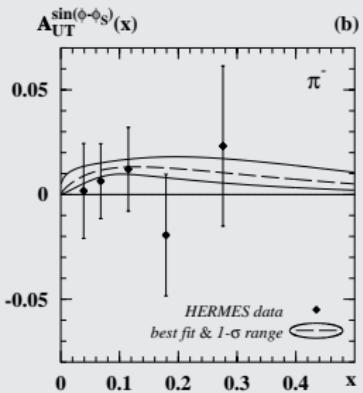
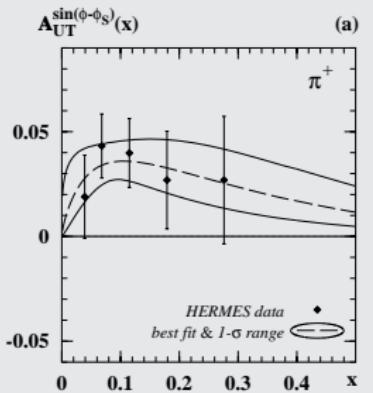
$$A_{UT}^{\sin(\phi_H - \phi_S)} \propto \frac{f_{1T}^\perp(x) D_1(z)}{f_1(x) D_1(z)}$$



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Hermes data



HERMES
Collaboration,
Phys.Rev.Lett. 94
(2005) 012002



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Sivers functions and its parametrizations

- we fit the first moment of the Sivers function
- $$f_{1T}^{\perp(1)}(x) = \int d\vec{k}_T \frac{\vec{k}_T^2}{2M_N^2} f_{1T}^{\perp}(x, \vec{k}_T^2)$$
- k_T -dependence is assumed to be Gaussian

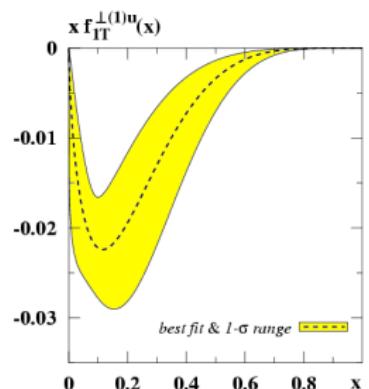
Fit to HERMES data

Reduce number of parameters

- $f_{1T}^{\perp, s, \bar{q}, g} = 0$

- Large - $N_C \Rightarrow f_{1T}^{\perp, u} = -f_{1T}^{\perp, d}$ [Pobylitsa, hep-ph/0212027]

⇒ only need to fit $f_{1T}^{\perp, u}$



- Ansatz for Sivers function

$$xf_{1T}^{\perp(1)} = Ax^B(1-x)^5$$

(also checked with $(1-x)^4$)

- Gaussian k_\perp and p_\perp dependence



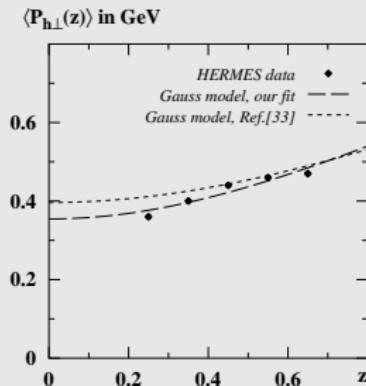
Phys.Rev. D73 (2006) 014021

Determining the gauss width

$$f_1(x, \vec{k}_\perp) = f_1(x) \exp \cdot \left(-\frac{\vec{k}_\perp^2}{k_{unp}^2} \right)$$

$$f_{1T}^\perp(x, \vec{k}_\perp) = f_{1T}^\perp(x) \exp \cdot \left(-\frac{\vec{k}_\perp^2}{k_{Siv}^2} \right)$$

$$D_1(z, \vec{p}_\perp) = D_1(z) \exp \cdot \left(-\frac{\vec{p}_\perp^2}{p_{D1}^2} \right)$$



$$\langle P_{hT}(z) \rangle_{Siv} \stackrel{\text{Gauss}}{=} \frac{2}{\sqrt{\pi}} \sqrt{z^2 k_{unp}^2 + p_{D1}^2}$$



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Determining the Gauss width

- \Rightarrow 3 Parameters to fit from external conditions
- p_{D1}^2, k_{unp}^2 from HERMES $P_{h\perp}$ -data

$$p_{D1}^2 = 0.16 \text{ GeV}^2 \quad k_{unp}^2 = 0.33 \text{ GeV}^2$$

- use positivity constraint

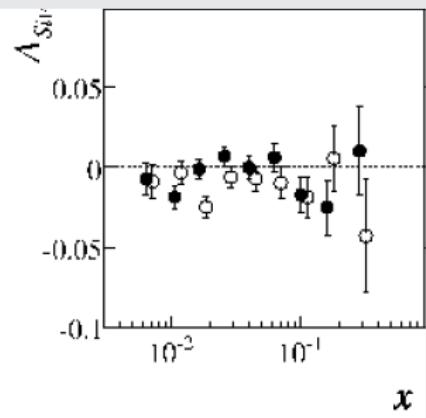
$$\frac{|p_\perp|}{M} f_{1T} \left(x, \vec{k}_\perp^2 \right) \leq f_1 \left(x, \vec{k}_\perp^2 \right) \Rightarrow k_{siv}^2 \leq \frac{k_{unp}^2}{1 + \frac{2M_N^2 k_{unp}^2}{ek_{siv}^2} \left(\frac{f_{1T}(x)}{f_1(x)} \right)^2}$$

- $\Rightarrow 0 < k_{siv}^2 < k_{unp}^2$
- k_{siv}^2 -range gives only 10% inaccuracy
- choose average $k_{siv}^2 = 0.15 \text{ GeV}^2$



Comparison to COMPASS data

COMPASS does not measure any significant Sivers asymmetry!



COMPASS uses a Deuterium target
⇒ number of u - and d -quarks is equal.
Large- N_C Ansatz implies this result naturally, because

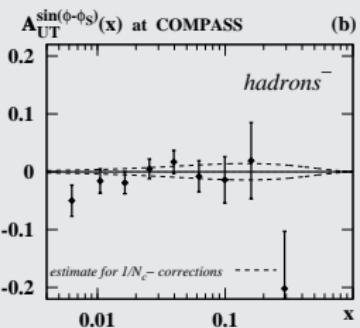
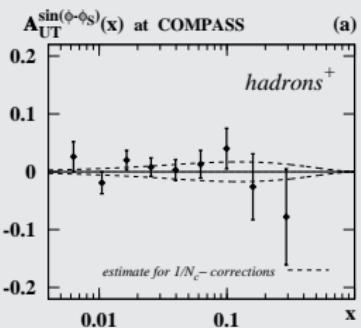
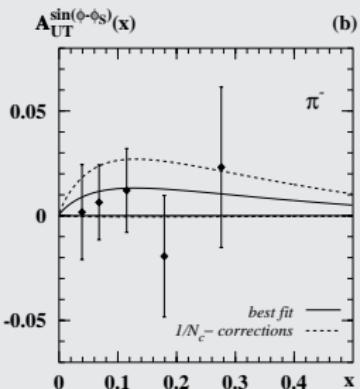
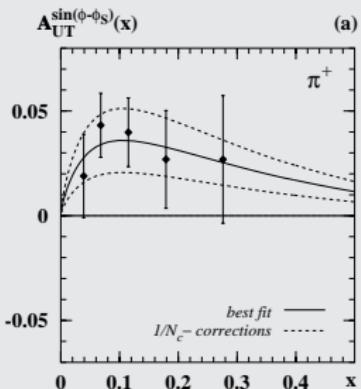
$$f_{1T}^{\perp,u} = -f_{1T}^{\perp,d}$$



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Validity of Large- N_C Ansatz



Comparison to other groups

hep[-ph/0511017]

Phys. Rev. D 71 074006 J.C.Collins, A.V.Efremov, K.Goeke,
S.Menzel, A.Metz, P.Schweitzer :

$$f_{1T}^{\perp(1),u} = Ax^b(1-x)^5 \frac{\pi}{\langle k_\perp^2 \rangle} e^{-\frac{\vec{k}_\perp^2}{\langle k_\perp^2 \rangle}}$$

Phys. Rev. D 70 117504 M.Anselmino, M.Boglione, U.D'Alesio,
A.Kotzinian, F.Murgia, A.Prokudin :

$$\Delta^N f_{q/p^\perp} = 2N_q(x) f_{q/p}(x) g\left(\vec{k}_\perp^2\right) h\left(\vec{k}_\perp^2\right)$$

$$N_q(x) = N_q x^{a_q} (1-x)^{b_q} \frac{(a_q+b_q)^{(a_q+b_q)}}{a_q^{a_q} b_q^{b_q}}$$

$g\left(\vec{k}_\perp^2\right)$, $h\left(\vec{k}_\perp^2\right)$ Gaussian or other function

Phys. Rev. D 72 054028 W.Vogelsang, F.Yuan:

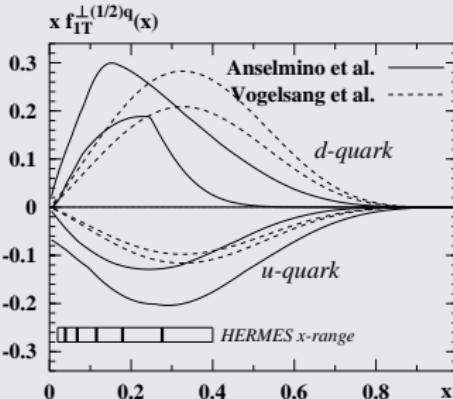
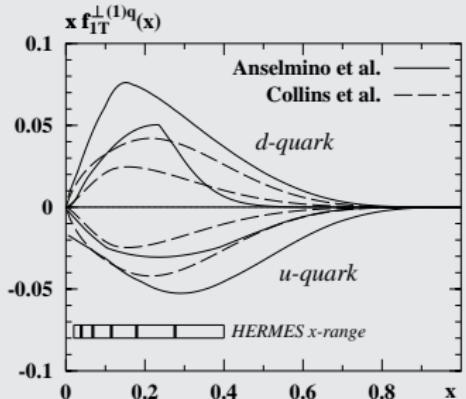
$$F_{siv}(x) = f_{1T}^{\perp(1/2)} = \frac{1}{2} \int d^2 \vec{k}_\perp \frac{|\vec{k}_\perp|}{M_N} f_{1T}^\perp(x, \vec{k}_\perp^2)$$

\vec{p}_\perp of D_1^q is neglected



Comparison to other groups

[hep-ph/0511017]



- all 3 describe the data equally well
- all 3 describe only $f_{1T}^{\perp u,d}$



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Gauge-Link

- ensures the color gauge-invariance of the quark operator
- generated by so-called final state interactions (FSI)
- the scattered parton can exchange gluons with the target remnant
- also called Wilson line

Definition

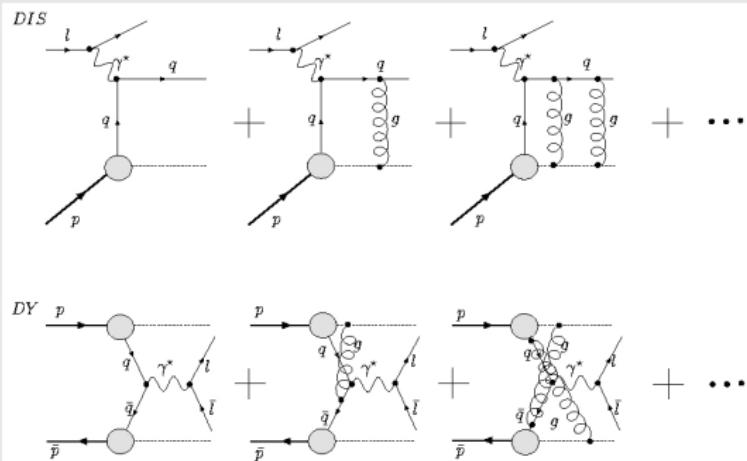
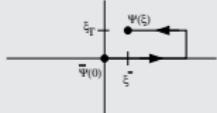
$$\mathcal{W} = \mathcal{P} \exp \left(ig \int d\alpha^\mu A_\mu(\alpha) \right)$$



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Gauge-Link in SIDIS and in Drell-Yan

 $W[0, \xi, \text{SIDIS}]$  $W[0, \xi, \text{DY}]$ 

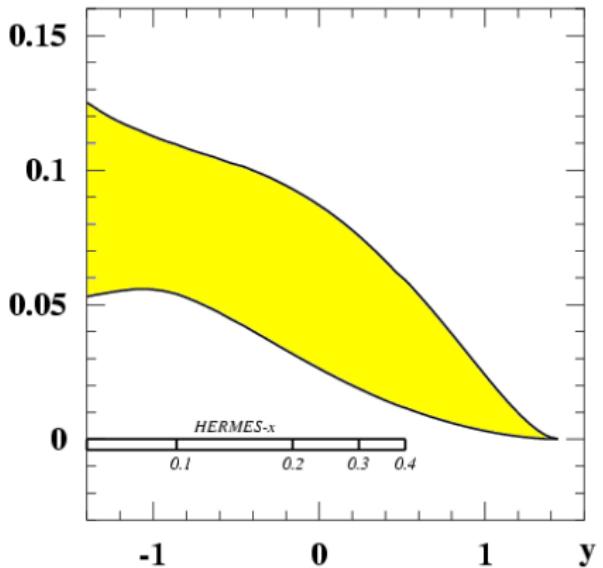
Gauge-Link

Gauge-Link in SIDIS and in DY run the other way round
Collins, Phys. Lett **B** 536, 43 (2002)



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DY@PAX $A_{UT}^{\sin(\phi - \phi_S)}$ in $p^+ \bar{p} \rightarrow l^+ l^- X$ at PAX

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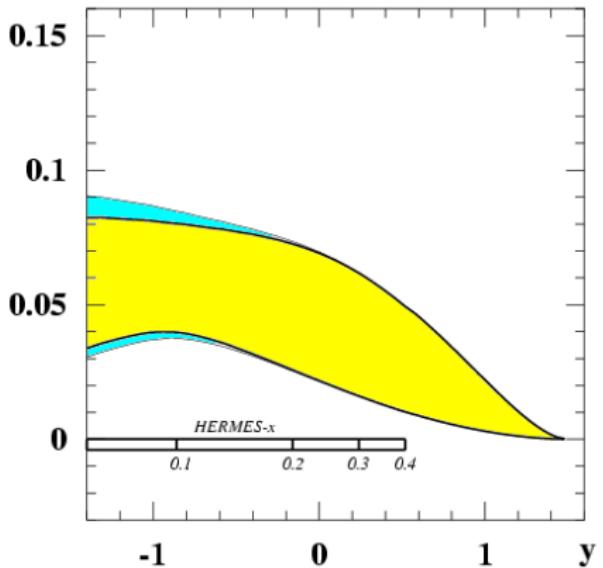
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Sivers functions and its parametrizations



DY@COMPASS

$A_{UT}^{\sin(\phi - \phi_S)}$ in $p^\uparrow \pi^- \rightarrow l^+ l^- X$ at COMPASS



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DY@RHIC → antiquark

\bar{q} can give significant contribution

PAX: Collision of protons and antiprotons

COMPASS: Collision of protons an π^-

⇒ Sivers antiquark distributions hardly visible

RHIC: Proton-proton collisions

sensitive to $f_{1T,DY}^{\perp,u}(x_1)f_1^{\bar{u}}(x_2)$ and $f_1^u(x_1)f_{1T,DY}^{\perp,\bar{u}}(x_2)$ on equal footing

⇒ Antiquarks should give significant contribution!

Models for Antiquark distributions

$$① \quad f_{1T}^{\perp\bar{q}} = 25 \% \ f_{1T}^{\perp q}$$

$$② \quad \frac{f_{1T}^{\perp,q}}{f_{1T}^{\perp,\bar{q}}} = \frac{f_1^u + f_1^d}{f_1^{\bar{u}} + f_1^{\bar{d}}}$$

Results

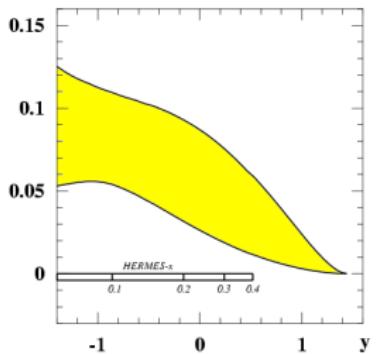
Measurement can give first boundaries on \bar{q} distributions



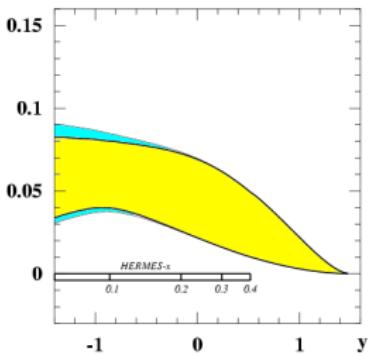
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Antiquarks at PAX and COMPASS

$A_{UT}^{\sin(\phi - \phi_S)}$ in $p^\dagger \bar{p} \rightarrow l^+ l^- X$ at PAX



$A_{UT}^{\sin(\phi - \phi_S)}$ in $p^\dagger \pi^- \rightarrow l^+ l^- X$ at COMPASS



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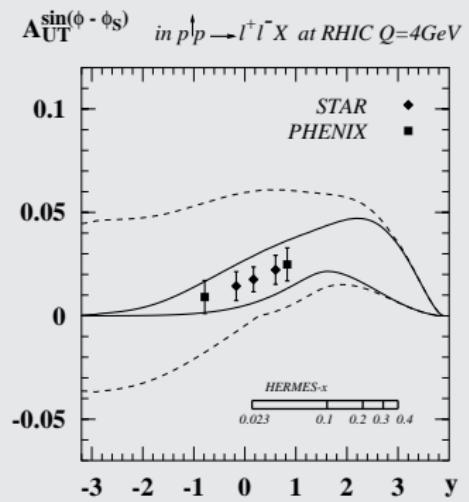
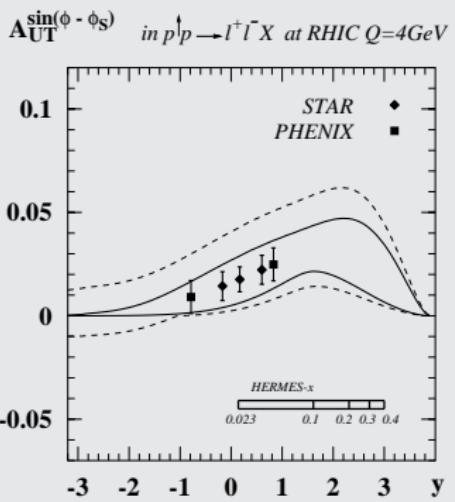
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DY@RHIC some $Q = 4 \text{ GeV}$

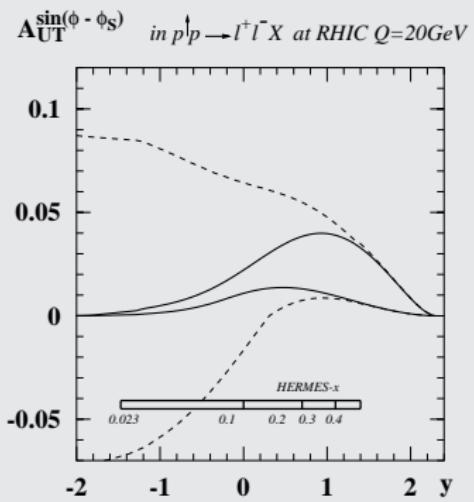
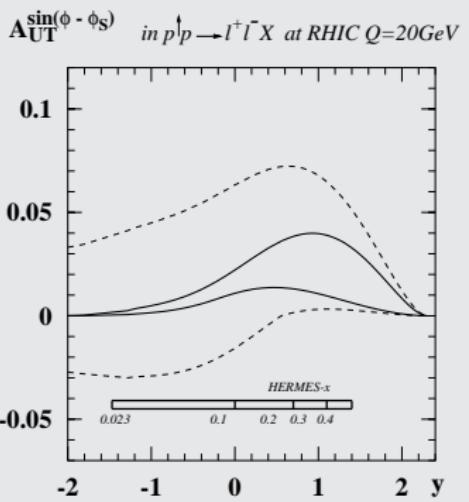
J.C. Collins, A.V. Efremov, K. Goeke, M. Grosse Perdekamp,
 S. Menzel, B. Meredith, A. Metz, P. Schweitzer
 [Phys.Rev. **D73** (2006) 094023]



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DY@RHIC $Q = 20 \text{ GeV}$

J.C. Collins, A.V. Efremov, K. Goeke, M. Grosse Perdekamp,
 S. Menzel, B. Meredith, A. Metz, P. Schweitzer
 [Phys.Rev. **D73** (2006) 094023]



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Summary

- Definition and descriptions of the Sivers function
 - Properties
 - Interpretation
- Parameterizations of the Sivers function
 - Gauss-Ansatz and Large N_C
 - Comparison of different parameterizations
- Future measurements of the Sivers function
 - Testing QCD (change of sign in SIDIS → DY)
 - Sivers antiquarks (RHIC)



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Outlook

- Get more accurate data of A_{UT}
 - Decide which parameterization is right
 - Relax Large- N_C ansatz
- Future experimental data
 - A_{UT}^{SIDIS} :
 - COMPASS (proton target)
 - HERMES (weighted data, more statistics)
 - hadronic final states
 - RHIC ($p^\dagger p \rightarrow jet_1 jet_2 X; jet H X; etc.$)
 - RHIC ($p^\dagger p$)
 - A_{UT}^{DY} :
 - PAX ($p^\dagger \bar{p}$)
 - COMPASS ($p^\dagger \pi^-$)
 - RHIC ($p^\dagger p$, Antiquarks)



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